

# The First Principle Study of AlX (X=4d elements) Dimer\*

## 铝-4d 过渡族金属二聚体的第一性原理计算

OUYANG Yi-fang, FANG Jie, BAO Jian-hui, WANG Rong-cheng, TAO Xiao-ma  
欧阳义芳, 方杰, 包建辉, 王戎丞, 陶小马

(College of Physical Science and Technology, Guangxi University, Nanning, Guangxi, 530004, China)

(广西大学物理科学与工程技术学院, 广西南宁 530004)

**Abstract:【Objective】**The nature of chemical bond between a transition metal and aluminium is a base of chemistry and materials science. The physical properties of AlX (X=4d elements) Dimer were studied. **【Methods】**Bond length, vibrational frequency, dissociation energy, ionization potential, energy gap and dipole moment of the title molecules in neutral and charged ions were studied by density functional method B3LYP. **【Results】**For neutral and charged molecules, the bond length decreases from AlY (AlY<sup>+</sup> and AlY<sup>-</sup>) to AlPd (AlPd<sup>+</sup> and AlPd<sup>-</sup>) except AlMo, AlNb<sup>+</sup> and AlMo<sup>+</sup>, and then increases. Opposite trends are observed for dissociation energy and vibrational frequency. The ionization energy through the series increases and reaches the maximum at AlRh and then decreases. The present results are in good agreement with the available experimental and other theoretical values except the vibrational frequency of AlTc. The present calculations show that the late transition metal can combine strongly with aluminum compared with the former transition metal (except AlCd). For cationic dimer, the dissociation procedure of channel AlX<sup>+</sup>→X + Al<sup>+</sup> is easier than that of channel AlX<sup>+</sup>→Al + X<sup>+</sup>. However, for anionic dimer, the channel AlX<sup>-</sup>→Al + X<sup>-</sup> is favorable except AlY<sup>-</sup>, AlPd<sup>-</sup> and AlCd<sup>-</sup>. The ground state spin multiplicity of the charged ions differs by ±1 from that of the corresponding neutral parent. The bond length of AlX<sup>+</sup> is larger than that of AlX (except AlPd) and that of AlX<sup>-</sup> (except AlRu<sup>-</sup> and AlPd<sup>-</sup>). The dissociation energy of charged dimer is bigger than that of neutral counterpart.

**Key words:** dimer, first-principles, transition metal aluminides, bonding characteristic

**摘要:【目的】**过渡族金属与铝之间的化学成键特性是化学与材料等领域的基础,对铝-4d 过渡族金属二聚体的物理性能进行研究。**【方法】**用密度泛函 B3LYP 方法计算铝-4d 过渡族金属二聚体的键长、振动频率、解离能、电离势、能隙和二极矩。**【结果】**对中性 and 荷电的二聚体,除 AlMo、AlNb<sup>+</sup> 和 AlMo<sup>+</sup> 外其键长从 AlY 到 AlPd 先减少后增加,解离能和振动频率的变化趋势与之相反。电离势随原子序数的递增先增加,在 AlRh 处达到极大值后减少。除 AlTc 的振动频率外,计算结果与已有的实验结果和其他的理论计算结果相符。计算结果还显示除 AlCd 外,前过渡族金属与铝的结合强度比后过渡族金属的结合要强。荷电二聚体的自旋多重度比相应的中性二聚体差±1。带正电的二聚体的键长比中性和带负电的二聚体的键长要长,荷电二聚体的解离能大于中性

二聚体的解离能,解离方式为 AlX<sup>+</sup>→X + Al<sup>+</sup> 和 AlX<sup>-</sup>→Al + X<sup>-</sup>。

**关键词:** 二聚体 第一性原理 过渡族金属铝化物 成键特征

**中图分类号:** TG1 **文献标识码:** A **文章编号:** 1005-9164(2014)03-0209-06

收稿日期:2014-05-16

作者简介:欧阳义芳(1965-),男,博士,教授,主要从事新材料制备和材料性能研究。

\* 国家自然科学基金(51061004,50831007)资助。

Understanding the nature of chemical bond between a transition metal and aluminium is of interest in many areas of science, e. g. surface chemistry, catalysis, astrophysics and organometallic chemistry, high temperature chemistry, materials science and astrochemistry. Diatomic metal aluminides are the simplest systems to understand the metal-metal bonding. A thorough understanding of how these species interact is an essential prerequisite to increase our knowledge of more complicated systems. Besides pure transition metal diatomic, the dimers containing transition metal have gotten more and more interests due to their application in catalysis, microelectronics and optical materials. Recently, the structures and stability of  $AlX$  and  $AlX^-$  species ( $X=H, Li, B, C, N, O$  and  $F$ ) have been reviewed<sup>[1]</sup>. For transition metal hydrides TMH ( $TM = Sc - Cu$ ), diverse studies were available. These include theoretical studies by density functional theory (DFT) of both neutral and cations<sup>[2]</sup>, modified coupled-pair functional (MCPF) on both of neutral<sup>[3]</sup> and cationic species<sup>[4]</sup>, pseudopotential multiconfiguration self-consistent field (MCSCF) on neutral species<sup>[5]</sup>, and local density approximation (LDA) on cations<sup>[6]</sup>. Some theoretical studies have been done for metal aluminides<sup>[7,8]</sup>. However, for transition metal aluminides  $AlTM$  ( $TM = 4d$  elements) and their charged ions, there is a little theoretical study for spectroscopic constants of the neutral dimers<sup>[8]</sup>. There is no experimental data have been reported on the spectroscopic constants of  $4d$   $AlX$  dimers except  $AlY$ <sup>[9]</sup>,  $AlPd$ <sup>[10,11]</sup> and  $AlAg$ <sup>[12~14]</sup>.

Density functional theory is currently widely used to determine structures of a wide variety of molecules. Compared to high-level *ab initio* molecular orbital theories, DFT has the advantage of applicability anywhere in periodic table and inherent computational efficiency. In this paper, we examined the  $4d$ -metal aluminides using DFT. The results are compared with the available theoretical and experimental values. It is hoped that our computational results will stimulate further experimental studies on these transition metal aluminides.

## 1 Computational methods

Bond length ( $r_e$ ), vibrational frequency ( $\omega$ ),

dissociation energy ( $D_0$ ), ionization energy ( $IE$ ) and electron affinity ( $EA$ ) for neutral and charged dimers (without  $IE$  and  $EA$ ) were determined for the  $AlX$ ,  $AlX^+$  and  $AlX^-$  ( $X=4d$  metals) using density functional method B3LYP, which includes a mixture of a traditional Hartree-Fock-like exchange energy, the Slater exchange functional, with gradient corrections due to Beck<sup>[15]</sup>, and the correlation potential of Vosko, Wilk and Nusair, with gradient correction due to Lee et al<sup>[16]</sup>. The calculations presented here were carried out using the GAUSSIAN 03 suit of program<sup>[17]</sup>. The basis set included; the basis set denoted as 6-311++G(2df) for Al, the ECP used was the ECP28MDF<sup>[18]</sup> and basis set CC-PVDZ-PP for  $4d$  transition metals from the Stuttgart Group (<http://www.theochem.uni-stuttgart.de/index.html>). To avoid trapping at local minima of the potential energy curve, a number of different initial geometries (bond lengths) are adopted. Due to the spin polarization, the molecules were calculated at various possible spin multiplicities in order to find the global minimum or ground state for each molecule (neutral and charged species). The calculated energies were corrected by the zero-point vibrational energies (except  $AlTc$  which contains radioactive element Tc, and the vibrational frequency is not available by the present method).

The zero-point vibrational energies ( $Z_{ALX}$ ) are estimated within the harmonic approximation. The dissociation energy was computed as the difference in the total energies and the isolated atoms, for neutral species, the dissociation energy is

$$D_0(AlX) = E_{tot}(Al) + E_{tot}(X) - E_{tot}(AlX) - Z_{ALX}. \quad (1)$$

Within the Born-Oppenheimer approximation, one could evaluate the  $EA$  and  $IE$

$$EA = E_{tot}(AlX) + Z_{AlX} - E_{tot}(AlX^-) - Z_{AlX^-}, \quad (2)$$

$$IE = E_{tot}(AlX^+) + Z_{AlX^+} - E_{tot}(AlX) - Z_{AlX}. \quad (3)$$

The energy gap  $E_G$  is

$$E_G = E_{Lumo} - E_{Homo}. \quad (4)$$

## 2 Results and discussion

The calculated spectroscopic constants of the

neutral 4d -transition metal aluminides were compiled in Table 1. The calculated spectroscopic constants of charged molecules are given in Tables 2 and 3.

## 2.1 Neutral dimer

AlY: It can be seen from Table 1 that  $^3\Sigma^-$  with electron configuration  $7\sigma^2 3\pi^2$  is the ground state, which is in agreement with experiment<sup>[9]</sup>. We also found that the bonding is probably described by valence electrons consisting of the 3p electron of aluminum and 4d electrons of yttrium, which combine to give a ground molecular configuration of  $\sigma^2\pi^2$ . It appears that the 3p orbits of aluminum are heavily involved with the 4d orbits. The calculated bond length 2.8429Å and dissociation energy 1.2647eV are close to experimental values 2.8728Å and 1.146eV, respectively.

AlZr:  $^6\Delta$  with electronic configuration  $7\sigma^1 8\sigma^1 3\pi^2 \delta^1$  is the ground state. Comparing the e-

lectronic configuration of AlZr with AlY, a dramatic bonding pattern change is observed. This is because the 5s orbital (main component of molecular orbital  $6\sigma$ ) of Y lies below 4d, whereas it is opposite for Zr<sup>[19]</sup>. The calculated dipole moment is the largest value in the neutral molecular AlX (X = 4d elements).

AlNb:  $^7\Sigma^+$  (with electron configuration  $7\sigma^1 8\sigma^1 3\pi^2 \delta^2$ ) is the ground state and AlNb possesses the largest spin polarization among the studied AlX species. From AlY, AlZr to AlNb, the bond length decreases, while the vibrational frequency, dissociation energy and ionization energy increase.

AlMo: The present calculation shows that  $^6\Sigma^+$  (with electron configuration  $8\sigma^1 3\pi^2 \delta^2$ ) is the ground state. The calculated bond length 2.7390Å of AlMo is only smaller than the values of AlY, but bigger than the values of AlZr and AlNb.

**Table 1 Physical parameters of the neutral 4d -transition metal aluminide AlX**

Neutral dimer	Basis set	State	Electronic configuration	$r_e$	$\omega$	$\mu$	$D_0$	IE	EA	$E_G$
AlY	CC-PVDZ-PP	$^3\Sigma^-$	$7\sigma^2 3\pi^2$	2.8429	247.1	1.7718	1.2647	5.7694	1.0276	1.7886
	Lan12dz <sup>[8]</sup>	$^3\Sigma^-$		2.8503	241.3		1.3253			
	Expt <sup>[9]</sup>	$^3\Sigma^-$		2.8728			1.146			
AlZr	CC-PVDZ-PP	$^6\Delta$	$7\sigma^1 8\sigma^1 3\pi^2 \delta^1$	2.6402	292.5	2.3987	1.5219	6.0024	1.2549	1.5551
	Lan12dz <sup>[8]</sup>	$^6\Delta$		2.6505	291.4		1.7139			
AlNb	CC-PVDZ-PP	$^7\Sigma^+$	$7\sigma^1 8\sigma^1 3\pi^2 \delta^2$	2.5825	296.2	0.9064	1.7964	6.0439	1.0741	1.5970
	Lan12dz <sup>[8]</sup>	$^7\Sigma^+$		2.5958	294.3		1.7881			
AlMo	CC-PVDZ-PP	$^6\Sigma^+$	$8\sigma^1 3\pi^2 \delta^2$	2.7390	226.0	0.5870	1.3919	6.1107	0.8461	1.8496
	Lan12dz <sup>[8]</sup>	$^6\Sigma^+$		2.7564	221.2		1.3519			
AlTc	CC-PVDZ-PP	$^5\Sigma^-$	$7\sigma^1 8\sigma^1 3\pi^4 \delta^2$	2.3119		0.2960	2.1012	6.5543	1.1823	1.8868
	Lan12dz <sup>[8]</sup>	$^5\Sigma^-$		2.3219			2.1339			
AlRu	CC-PVDZ-PP	$^4\Sigma^-$	$8\sigma^1 3\pi^4 \delta^2$	2.2628	342.8	0.5616	2.4146	6.9475	0.8696	1.8994
	Lan12dz <sup>[8]</sup>	$^4\Sigma^-$		2.2738	357.7	2.4054				
AlRh	CC-PVDZ-PP	$^3\Delta$	$8\sigma^1 3\pi^4 \delta^3$	2.2359	377.1	0.1883	2.8169	7.4731	1.1162	1.8240
	Lan12dz <sup>[8]</sup>	$^3\Delta$		2.2551	372.9		2.8394			
AlPd	CC-PVDZ-PP	$^2\Sigma^+$	$8\sigma^1 3\pi^4 \delta^4$	2.2323	365.9	0.6511	2.7341	6.6699	1.0750	1.6969
	Lan12dz <sup>[8]</sup>	$^2\Sigma^+$		2.2509	361.5		2.6771			
	Expt <sup>[11]</sup>						2.591			
AlAg	CC-PVDZ-PP	$^1\Sigma^+$	$8\sigma^2 3\pi^4 \delta^4$	2.5112	248.4	0.8698	1.990	6.9335	0.6519	2.5954
	Lan12dz <sup>[8]</sup>	$^1\Sigma^+$		2.5515	237.2		1.8890			
	Expt <sup>[12]</sup>	$^1\Sigma^+$		2.474	256.1		2.2			
Expt <sup>[13]</sup>	$^1\Sigma^+$					1.95				
	ICMRCI+Qe <sup>[14]</sup>	$^1\Sigma^+$		2.5379	240					
AlCd	CC-PVDZ-PP	$^2\Pi$	$8\sigma^2 4\pi^1 \delta^4$	2.9157	125.7	0.5504	0.2137	5.6330	0.8236	1.3494
	Lan12dz <sup>[8]</sup>	$^2\Pi$		3.0794	108.7	0.1700				

Bond length ( $r_e$  in Å), vibrational frequency ( $\omega$  in  $\text{cm}^{-1}$ ), dipole moment ( $\mu$  in Debye), dissociation energy ( $D_0$  in eV), ionization energy (IE in eV), electron affinity (EA in eV) and energy gaps ( $E_G$  in eV).

AlX (X=Tc, Ru, Rh): No experimental data are available for these dimers. Our calculations indicate that the ground state spin multiplicities are 5 (with electron configuration  $7\sigma^1 8\sigma^1 3\pi^4 \delta^2$ ) for AlTc, 4 (electron configuration  $8\sigma^1 3\pi^4 \delta^2$ ) for AlRu, 3 (electron configuration  $8\sigma^1 3\pi^4 \delta^3$ ) for AlRh. For AlTc, vibrational frequency cannot be obtained by the present method. The bond length decreased from AlTc to AlRh, but the dissociation energy and ionization energy increase. Although the ionization energy 6.9475 eV of AlRu is smaller than the value 7.4731 eV of AlRh, the electron affinity 0.8696 eV of AlRu is also smaller than the value 1.1162 eV of AlRh. The dissociation energy 2.8169 eV and ionization energy 7.4731 eV of AlRh are the biggest values, while the dipole moment 0.1883 Debye of AlRh is the smallest in the AlX.

AlPd:  $^2\Sigma^+$  with electron configuration  $8\sigma^1 3\pi^4 \delta^4$  is the ground state. The bond length 2.2323 Å is smallest in AlX, and our calculated dissociation energy 2.7341 eV of AlPd agrees with experimental value 2.591 eV<sup>[11]</sup>.

AlAg: There are several experimental and theoretical studies available for neutral dimer AlAg<sup>[12~14]</sup>. Our calculation shows that single state ( $^1\Sigma^+$ ) with electron configuration  $8\sigma^2 3\pi^4 \delta^4$  is the ground state. The calculated bond length 2.5112 Å, and vibrational frequency 248.4 cm<sup>-1</sup> are in good agreement with the previous experimental values 2.474 Å and 256.1 cm<sup>-1</sup><sup>[12]</sup>. The calculated dissociation energy (1.990 eV) agrees well with experimental value 1.95 eV<sup>[13]</sup>. However, our dissociation energy underestimates experimental value 2.2 eV<sup>[12]</sup> by 0.21 eV. The ICMRCI+Q results<sup>[14]</sup> (2.5379 Å and 240 cm<sup>-1</sup>) for the  $^1\Sigma^+$  state of AlAg are in excellent agreement with our results. For the noble metals, the electronic structure of silver is best approximated by the closed shell  $d^{10} s^1$  configuration.  $^1\Sigma^+$  results from a single bond between the Ag 5s and Al 3p valence electrons.

AlCd: For AlCd,  $^2\Pi$  with electron configuration  $8\sigma^2 4\pi^1 \delta^4$  is the ground state. As closed shell of cadmium ( $(4d^{10} 5s^2)$ ), the bond length 2.9157 Å is the largest value, and the vibrational frequency 125.7 cm<sup>-1</sup>, dissociation energy 0.2137 eV, ionization

energy 5.6330 eV and energy gap 1.3494 eV are the smallest in AlX, which indicated relatively weak bond.

In a whole, the results calculated in the present are in good agreement with those obtained by B3LYP with LanL2dz for 4d-metals<sup>[8]</sup>.

## 2.2 Cationic dimer

The electron configurations of ground state for AlX<sup>+</sup> are given in Table 2. There is a recent ESR (Electron spin resonance) study<sup>[20]</sup> of AlY<sup>+</sup> indicating that both molecules possess  $^4\Sigma^-$  ground states deriving from  $\sigma^1 \pi^2$  configurations. The difference between the present calculation and experiment maybe result from the basis set for the 4d metal used in the present calculation. The calculation was performed by using LanL2dz instead of ECP28MDF for Y. The results indicate that the ground state is  $^4\Sigma^-$ . Further experiment is needed to verify it. One can see that the multiplicity of ground state AlX<sup>+</sup> differs from the counterpart of neutral one by one. The similar phenomenon was also observed for Al-3d transition metal dimers.

The bond length of AlX<sup>+</sup> is larger than that of corresponding AlX except AlRh and AlPd, for which the opposite phenomenon appears. The vibrational frequency of AlX<sup>+</sup> is lower than that of corresponding AlX. The bond length of AlPd<sup>+</sup> is the shortest in the series of AlX<sup>+</sup>, while its vibrational frequency is the largest. The bond length of AlZn<sup>+</sup> is the largest and the vibrational frequency is the lowest.

The dipole moment of AlX<sup>+</sup> is bigger than that of neutral counterpart except AlY and AlZr. As for the energy gap between the LUMO and HOMO levels, the values of AlX<sup>+</sup> are bigger than those of AlX except AlRh and AlAg.

From Table 2, one can see that the calculated dissociation energies indicate that the dissociation channel of AlX<sup>+</sup> → X + Al<sup>+</sup> is favorable for AlX<sup>+</sup>. The dissociation energies of AlX<sup>+</sup> are smaller than those of AlX except AlY, AlZr and AlCd. This indicates that neutral dimer is more stable than cationic dimer for most of Al-4d transition metals.

## 2.3 Anionic dimer

In Table 3, the electronic configurations for AlX<sup>-</sup> are compiled. It can be seen that the spin multiplicity of ground state AlX<sup>-</sup> differs from the corresponding AlX one by one, which is similar as that for AlX<sup>+</sup> mentioned above.

**Table 2** Physical parameters of the cation  $4d$ -transition metal aluminide  $\text{AlX}^+$ 

Cation dimer	State	Electronic configuration	$r_e$	$\omega$	$\mu$	$D_0$ $\text{AlX}^+ \rightarrow$ $\text{X}^+ + \text{Al}$	$D_0$ $\text{AlX}^+ \rightarrow$ $\text{X} + \text{Al}^+$	$E_G$
$\text{AlY}^+$	$^2\Pi$	$7\sigma^2 3\pi^1$	3.2379	148.6	0.6994	1.8344	1.5165	1.7296
$\text{AlZr}^+$	$^5\Sigma^-$	$7\sigma^1 8\sigma^1 3\pi^2$	2.7347	241.5	0.2717	2.2013	1.5408	1.7560
$\text{AlNb}^+$	$^6\Delta$	$7\sigma^1 8\sigma^1 3\pi^2 \delta^1$	2.7434	221.0	1.8111	2.9577	1.7737	2.0520
$\text{AlMo}^+$	$^7\Sigma^+$	$7\sigma^1 8\sigma^1 3\pi^2 \delta^2$	2.7835	191.5	3.0695	2.5461	1.3020	2.2863
$\text{AlTc}^+$	$^6\Phi$	$7\sigma^1 8\sigma^1 9\sigma^1 3\pi^3 \delta^1$	2.5874		2.9417	2.6586	1.5682	2.0743
$\text{AlRu}^+$	$^3\Sigma^-$	$7\sigma^2 3\pi^4 \delta^2$	2.3389	283.6	2.4647	3.0975	1.4883	2.0972
$\text{AlRh}^+$	$^2\Sigma^+$	$7\sigma^2 3\pi^4 \delta^4$	2.2573	351.9	4.5018	3.0850	1.6989	1.4569
$\text{AlPd}^+$	$^1\Sigma^+$	$7\sigma^2 3\pi^4 \delta^4$	2.2265	349.3	4.2999	4.6521	2.0854	3.5843
$\text{AlAg}^+$	$^2\Sigma^+$	$8\sigma^1 3\pi^4 \delta^4$	2.5774	188.9	4.4670	3.0012	1.0777	1.8909
$\text{AlCd}^+$	$^1\Sigma^+$	$8\sigma^2 3\pi^4 \delta^4$	3.1107	127.0	6.2584	3.8400	0.6019	3.8839

Bond length ( $r_e$  in Å), vibrational frequency ( $\omega$  in  $\text{cm}^{-1}$ ), dipole moment ( $\mu$  in Debye), dissociation energy ( $D_0$  in eV) and energy gaps ( $E_G$  in eV)

**Table 3** Physical parameters of the anion  $4d$ -transition metal aluminide  $\text{AlX}^-$ 

Anion dimer	State	Electronic configuration	$r_e$	$\omega$	$\mu$	$D_0$ $\text{AlX}^- \rightarrow$ $\text{X}^- + \text{Al}$	$D_0$ $\text{AlX}^- \rightarrow$ $\text{X} + \text{Al}^-$	$E_G$
$\text{AlY}^-$	$^4\Sigma^-$	$8\sigma^1 3\pi^2$	2.8378	252.3	4.3267	2.5801	1.9083	1.7418
$\text{AlZr}^-$	$^5\Delta$	$8\sigma^1 3\pi^2 \delta^1$	2.7156	267.8	2.6325	2.3139	2.3925	2.0104
$\text{AlNb}^-$	$^6\Sigma^+$	$8\sigma^1 3\pi^2 \delta^2$	2.6300	288.4	2.6769	2.2586	2.4865	1.7734
$\text{AlMo}^-$	$^5\Phi$	$8\sigma^1 9\sigma^1 3\pi^3 \delta^1$	2.6247	224.4	4.1319	1.7646	1.8536	1.2574
$\text{AlTc}^-$	$^4\Sigma^-$	$8\sigma^1 3\pi^4 \delta^2$	2.3942		3.7010	2.6016	2.8993	1.9668
$\text{AlRu}^-$	$^3\Delta$	$8\sigma^1 3\pi^4 \delta^3$	2.3662	323.7	3.6341	2.4812	2.9003	1.9886
$\text{AlRh}^-$	$^2\Delta$	$8\sigma^2 3\pi^4 \delta^3$	2.2508	389.1	2.5319	3.0427	3.5258	1.9132
$\text{AlPd}^-$	$^1\Sigma^+$	$8\sigma^2 3\pi^4 \delta^4$	2.2419	376.6	2.5660	3.5157	3.4252	1.8721
$\text{AlAg}^-$	$^2\Pi$	$8\sigma^2 4\pi^1 \delta^4$	2.5536	210.4	4.5684	1.6113	2.2580	1.1627
$\text{AlCd}^-$	$^3\Sigma^-$	$8\sigma^2 4\pi^2 \delta^4$	2.7434	166.5	5.8880	1.9685	0.6534	1.2338

Bond length ( $r_e$  in Å), vibrational frequency ( $\omega$  in  $\text{cm}^{-1}$ ), dipole moment ( $\mu$  in Debye), dissociation energy ( $D_0$  in eV) and energy gaps ( $E_G$  in eV)

The bond lengths of  $\text{AlX}^-$  are shorter than those of  $\text{AlX}^+$ . By comparing the bond lengths of  $\text{AlX}^-$  and  $\text{AlX}$ , the bond length of anionic dimer is longer for  $\text{AlZr}$ ,  $\text{AlTc}$ ,  $\text{AlRu}$ ,  $\text{AlPd}$  and  $\text{AlAg}$ . The vibrational frequencies of  $\text{AlX}^-$  are higher than that of corresponding  $\text{AlX}^+$ , while only the vibrational frequencies of  $\text{AlY}$ ,  $\text{AlRu}$  and  $\text{AlCd}$  are lower than those of the corresponding anionic dimer.

The dipole moments of  $\text{AlX}^-$  are larger than those of corresponding  $\text{AlX}$ , while they are less than those of  $\text{AlX}^+$  except  $\text{AlRh}^+$ ,  $\text{AlPd}^+$  and  $\text{AlCd}^+$ . The energy gap between the LUMO and HOMO levels for  $\text{AlX}^-$  is larger than that of corresponding  $\text{AlX}$  except for  $\text{AlY}$ ,  $\text{AlMo}$ ,  $\text{AlAg}$  and  $\text{AlCd}$ , while it is less than that of corresponding  $\text{AlX}^+$  except  $\text{AlY}^+$  and  $\text{AlZr}^+$ .

The dissociation procedure of  $\text{AlX}^-$  is different from that of  $\text{AlX}^+$ . As for  $\text{AlX}^-$  ( $\text{X} = \text{Y}, \text{Nb}, \text{Pd}$  and  $\text{Cd}$ ), the dissociation energy of channel  $\text{AlX}^- \rightarrow \text{X}^- + \text{Al}$  is less than that of channel  $\text{AlX}^- \rightarrow \text{X} + \text{Al}^-$ . Opposite way was observed for the rest of

$\text{AlX}^-$ . The dissociation energy for  $\text{AlX}^-$  is for larger than that of neutral dimer and cathionic dimer except  $\text{AlAg}$ , while the dissociation energy of anionic  $\text{AlAg}$  is less than that of neutral  $\text{AlAg}$ . The dissociation energy of  $\text{AlCd}^-$  is the least in series of  $\text{AlX}^-$ .

### 3 Conclusions

Density functional method B3LYP was used to study the spectroscopic constants of  $4d$ -metal aluminides dimers and charged ones. The electronic configuration, bond length, harmonic vibrational frequency, and dissociation energy have been calculated for neutral and charged dimers. The present calculations are in good agreement with the available experimental and other theoretical results. The ground state spin multiplicity of the charged ions differs by  $\pm 1$  from that of the corresponding neutral parent. For neutral and charged molecular, the late transition metals (except  $\text{Ag}$  and  $\text{Cd}$ ) are found to form stronger bonding with aluminum than former tran-

sition metals. The dissociation channel of  $\text{AlX}^+ \rightarrow \text{X} + \text{Al}^+$  is favor for  $\text{AlX}^+$ . However, for anionic dimer, the channel  $\text{AlX}^- \rightarrow \text{Al} + \text{X}^-$  is favorable except  $\text{AlY}^-$ ,  $\text{AlPd}^-$  and  $\text{AlCd}^-$ . The bond length of  $\text{AlX}^+$  is larger than that of  $\text{AlX}$  (except  $\text{AlPd}$ ) and that of  $\text{AlX}^-$  (except  $\text{AlRu}^-$  and  $\text{AlPd}^-$ ). The dissociation energy of charged dimer is bigger than that of neutral counterpart.

#### References:

[1] Gutsev G L, Jena P, Bartlett R J. Structure and stability of  $\text{AlX}$  and  $\text{AlX}^-$  species[J]. *J Chem Phys*, 1999, 110(6):2928-2935.

[2] Barone V, Adamo C. First-row transition-metal hydrides: a challenging playground for new theoretical approaches[J]. *Int J Quantum Chem*, 1997, 61(3):443-451.

[3] Chong D P, Langhoff S R, Bauschlicher Jr C W, et al. Theoretical dipole moments for the first-row transition metal hydrides[J]. *J Chem Phys*, 1986, 85(5):2850-2860.

[4] Pettersson L G M, Bauschlicher Jr C W, Langhoff S P, et al. Positive ions of the first- and second-row transition metal hydrides[J]. *J Chem Phys*, 1987, 87(1):481-492.

[5] Das G. A pseudopotential study of the iron-series transition metal hydrides[J]. *J Chem Phys*, 1981, 74(10):5766-5774.

[6] Ziegler T, Li j. Bond energies for cationic bare metal hydrides of the first transition series: a challenge to density functional theory[J]. *Can J Chem*, 1994, 72(3):783-789.

[7] Wang J C, Zhai D M, Guo F, et al. Theoretical study of spectroscopic parameters of alkali-Al and alkaline earth-Al dimers[J]. *Theoretical Chemistry Accounts*, 2008, 121(3):165-172.

[8] Ouyang Y F, Wang J C, Hou Y H, et al. First principle study of  $\text{AlX}$  ( $X = 3d, 4d, 5d$  elements and Lu) dimers[J]. *J Chem Phys*, 2008, 128(7):074305-074306.

[9] Arrington C A, Langenber J D, Pineger J C, et al. Spectroscopy of Jet-cooled  $\text{AlY}$ [J]. *J Phys Chem*, 1995, 99

(9):2589-2593.

[10] Knight Jr L B, Babb R, Hill D W, et al. Laser vaporization generation of the diatomic radicals  $\text{PdB}$ ,  $^{105}\text{PdB}$ ,  $\text{PdAl}$ ; electron spin resonance investigation in neon matrices at 4k[J]. *J Chem Phys*, 1992, 97(5):2987-2995.

[11] Coke D L, Gingerich K A, Chang C A. Atomization energies of gaseous  $\text{AlPd}$  and  $\text{Al}_2\text{Pd}$ [J]. *J Chem Soc, Faraday Trans 1*, 1976, 72:268-272.

[12] Robbins D L, Yeh C S, Pilgrim J S, et al. Photoionization electronic spectroscopy of  $\text{AlAg}$ [J]. *J Chem Phys*, 1994, 100(7):4775-4783.

[13] Herzberg G. *Molecular spectra and molecular structure I spectra of diatomic molecules*[M]. New York: Van Nostrand Reinhold, 1950.

[14] Bauschlicher C W, Langhoff Jr S R, Partridge H. The low-lying states of  $\text{AlCu}$  and  $\text{AlAg}$ [J]. *J Chem Phys*, 1994, 100(2):1219-1225.

[15] Becke A D. Density-functional thermochemistry III. The role of exact exchange[J]. *J Chem Phys*, 1993, 98(7):5648-5652.

[16] Lee C, Yang W, Parr R. Development of the Coll-Salvetti correlation-energy formula into a functional of the electron density[J]. *Phys Rev B*, 1988, 37(2):785-789.

[17] Frisch M J, Trucks G W, Schlegel H B, et al. *Gaussian 03*[S]. Gaussian, Inc, Pittsburgh, PA, 2003.

[18] Peterson K A, Figgen D, Dolg M, et al. Energy-consistent relativistic pseudopotentials and correlation consistent basis sets for the  $4d$  elements Y-Pd[J]. *J Chem Phys*, 2007, 126(12):124101-12.

[19] Hay P J, Wadt W R. Ab initio effective core potentials for molecular calculations. Potentials for main group elements Na to Bi[J]. *J Chem Phys*, 1985(1):284-292.

[20] Knight L B, Babb Jr R M, King G M, et al. Laser vaporization generation of  $\text{Y}^{10}\text{B}^+$ ,  $\text{Y}^{11}\text{B}^+$ , and  $\text{YAl}^+$  for electron spin resonance studies in neon matrices at 4k: Comparison with theoretical calculations [J]. *J Chem Phys*, 1993, 98:4404-4412.

(责任编辑:陈小玲)