

一类四阶粘性扩散方程广义解的唯一性*

Uniqueness of Solutions to a Fourth Order Viscous Diffusion Equation

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摘要: 基于椭圆算子, 证明初边值问题: $\frac{\partial u}{\partial t} - \lambda \frac{\partial}{\partial t} (\frac{\partial^2 u}{\partial x^2}) + \frac{\partial^4 \Phi(u)}{\partial x^4} = 0, (x, t) \in Q_T, u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, t \in (0, T), u(x, 0) = u_0(x), x \in (0, 1), \lambda \geq 0$ 是粘性系数, $Q_T = (0, 1) \times (0, T), \Phi(u) = |u|^{q-2}u, q > 1$, 最多存在一个 L^2 解.

关键词: 扩散方程 粘性 解 唯一性 椭圆算子

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Abstract: Based on elliptic operators, the initial boundary value problem: $\frac{\partial u}{\partial t} - \lambda \frac{\partial}{\partial t} (\frac{\partial^2 u}{\partial x^2}) + \frac{\partial^4 \Phi(u)}{\partial x^4} = 0, (x, t) \in Q_T, u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, t \in (0, T), u(x, 0) = u_0(x), x \in (0, 1), \lambda \geq 0, Q_T = (0, 1) \times (0, T), \Phi(u) = |u|^{q-2}u, q > 1$, has at most one L^2 solution.

Key words: diffusion equation, viscous, solution, uniqueness, elliptic operator

考虑初边值问题:

$$\frac{\partial u}{\partial t} - \lambda \frac{\partial}{\partial t} (\frac{\partial^2 u}{\partial x^2}) + \frac{\partial^4 \Phi(u)}{\partial x^4} = 0, (x, t) \in Q_T, \quad (1)$$

$$u(0, t) = u(1, t) = u_x(0, t) = u_x(1, t) = 0, t \in (0, T), \quad (2)$$

$$u(x, 0) = u_0(x), x \in (0, 1), \quad (3)$$

其中 $\lambda \geq 0$ 是粘性系数, $Q_T = (0, 1) \times (0, T), \Phi(u) = |u|^{q-2}u, q > 1$.

问题(1) 是自然界中广泛存在的某些现象的数学模型, 当 $\lambda = 0$ 时, 得到一类四阶非线性扩散方程:

$$\frac{\partial u}{\partial t} + \frac{\partial^4 \Phi(u)}{\partial x^4} = 0.$$

此方程在满足 $\Phi'(s) = 0$ 的点是退化的, 而且解的

存在性与唯一性问题已经解决^[1]. 当 $\lambda > 0$ 时, 文献[2] 研究了二阶扩散方程:

$$\frac{\partial u}{\partial t} - \lambda \frac{\partial \Delta u}{\partial t} = \Delta A(u) + \text{div} \vec{B}(u),$$

也证明了其广义解的唯一性. 对于初边值问题(1) ~ (3), 文献[1] 讨论了其广义解的存在性, 文献[3] 讨论了当 $\lambda = 1$ 的情况, 并使用椭圆算子的方法, 证明了 L^2 解的唯一性. 受文献[3] 的启发, 本文讨论当 $\lambda \geq 0$ 时, 初边值问题(1) ~ (3) L^2 解的唯一性. 所得结果补充了文献[1] 中广义解唯一性的结论.

1 主要结果

定义 1 函数 $u(x, t)$ 称为初边值问题(1) ~ (3) 的 L^2 解, 如果 $u \in L^2(Q_T), \Phi(u) \in L^2(Q_T)$, 而且对于任意的检验函数 $\varphi(x, t) \in C^\infty(Q_T), \varphi(0, t) = \varphi(1, t) = \varphi_x(0, t) = \varphi_x(1, t) = \varphi_x(x, T) = 0$, 有积分等式

$$-\int_0^1 u_0(x) \varphi(x, 0) dx + \int_0^1 \lambda u_0(x) \frac{\partial^2 \varphi(x, 0)}{\partial x^2} dx -$$

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$$\iint_{Q_T} u(I - \lambda\Delta) \frac{\partial \varphi}{\partial t} dxdt + \iint_{Q_T} \Phi(u) \frac{\partial^4 \varphi}{\partial x^4} dxdt = 0 \quad (4)$$

成立,其中 $\Delta = \frac{\partial^2}{\partial x^2}$.

由定义1可以知道,对任意的函数 $\varphi \in H_0^2(Q_T) \cap H^4(Q_T)$,积分等式(4)也是成立的.

定理1 初边值问题(1)~(3)最多存在一个 L^2 解.

2 定理1的证明

假设 $u_1, u_2 \in L^2(Q_T)$ 是初边值问题(1)~(3)的两个 L^2 解.令 $z = u_1 - u_2, v = \Phi(u_1) - \Phi(u_2)$,根据 L^2 解的定义,可以得到

$$-\iint_{Q_T} z(I - \lambda\Delta) \frac{\partial \varphi}{\partial t} dxdt + \iint_{Q_T} v \frac{\partial^4 \varphi}{\partial x^4} dxdt = 0. \quad (5)$$

我们希望通过选择特殊的检验函数 φ 来证明(5)式在几乎处处的 Q_T 满足 $z = 0$.

对充分小的 $\mu > 0$,定义算子

$$T_\mu: L^2(I) \rightarrow H_0^2(I) \cap H^4(I), g \rightarrow y,$$

其中 $I = (0, 1), y = T_\mu g$ 是由两点边值问题

$$\frac{\partial^4 y}{\partial x^4} + \mu y = g, (\mu > 0),$$

$$y(0) = y(1) = y'(0) = y'(1) = 0,$$

唯一决定的.而且容易看出算子 T_μ 是自共轭的,也就是对任意的 $f, g \in L^2(I)$ 有

$$\int_0^1 g(T_\mu f) dx = \int_0^1 f(T_\mu g) dx.$$

由椭圆方程的 L^2 理论^[4],有 $T_\mu g \in H_0^2(I) \cap H^4(I)$ 满足

$$\begin{aligned} & \mu^2 \int_0^1 (T_\mu g)^2 dx + 2\mu \int_0^1 \left(\frac{\partial^2 T_\mu g}{\partial x^2}\right)^2 dx + \\ & \int_0^1 \left(\frac{\partial^4 T_\mu g}{\partial x^4}\right)^2 dx \leq \int_0^1 g^2 dx. \end{aligned} \quad (6)$$

注意如下的事实

$$\iint_{Q_T} z(I - \lambda\Delta) \frac{\partial \varphi}{\partial t} dxdt = \iint_{Q_T} (I - \lambda\Delta) z \frac{\partial \varphi}{\partial t} dxdt,$$

对任意的 $k(x) \in C_0^\infty(I), \Psi(t) \in C_0^\infty(0, T)$,在(5)式中用 $(T_\mu k)\Psi(t)$ 代替检验函数 φ ,得到

$$\begin{aligned} & -\iint_{Q_T} z(I - \lambda\Delta)(T_\mu k)\Psi'(t) dxdt + \\ & \iint_{Q_T} v \frac{\partial^4 T_\mu k}{\partial x^4} \Psi(t) dxdt = -\iint_{Q_T} (I - \\ & \lambda\Delta)(T_\mu z)\Psi'(t)k(x) dxdt + \iint_{Q_T} v(k - \\ & \mu T_\mu k)\Psi(t) dxdt = -\iint_{Q_T} (I - \end{aligned}$$

$$\lambda\Delta)(T_\mu z)\Psi'(t)k(x) dxdt + \iint_{Q_T} (v - \mu T_\mu v)\Psi(t)k(x) dxdt = 0.$$

由此知道,对任意的 $\varphi \in C_0^\infty(Q_T)$,有

$$\begin{aligned} & -\iint_{Q_T} (I - \lambda\Delta)(T_\mu z) \frac{\partial \varphi}{\partial t} dxdt + \iint_{Q_T} (v - \\ & \mu T_\mu v)\varphi dxdt = \iint_{Q_T} (I - \lambda\Delta) \frac{\partial T_\mu z}{\partial t} \varphi dxdt + \iint_{Q_T} (v - \\ & \mu T_\mu v)\varphi dxdt = 0. \end{aligned}$$

表明 $(I - \lambda\Delta) \frac{\partial T_\mu z}{\partial t} \in L^2(Q_T)$,并且在分布的意义下有

$$(I - \lambda\Delta) \frac{\partial T_\mu z}{\partial t} = \mu T_\mu v - v. \quad (7)$$

记

$$g_\mu(t) = \int_0^1 (I - \lambda\Delta) T_\mu z(x, t) \cdot z(x, t) dx$$

和证明

$$\lim_{\mu \rightarrow 0} g_\mu(t) = 0, a. e. t \in (0, T). \quad (8)$$

由于

$$\begin{aligned} & g_\mu(t) = \int_0^1 (I - \lambda\Delta) T_\mu z(x, t) \cdot z(x, t) dx = \\ & \int_0^1 (I - \lambda\Delta)(T_\mu z) \left(\frac{\partial^4 T_\mu z}{\partial x^4} + \mu T_\mu z\right) dx = \\ & \mu \int_0^1 (T_\mu z)^2 dx + \int_0^1 \left(\frac{\partial^2 T_\mu z}{\partial x^2}\right)^2 dx + \lambda \int_0^1 \left(\frac{\partial^3 T_\mu z}{\partial x^3}\right)^2 dx + \\ & \lambda \mu \int_0^1 \left(\frac{\partial T_\mu z}{\partial x}\right)^2 dx = \mu \|T_\mu z\|^2 + \left\| \frac{\partial^2 T_\mu z}{\partial x^2} \right\|^2 + \\ & \lambda \left\| \frac{\partial^3 T_\mu z}{\partial x^3} \right\|^2 + \lambda \mu \left\| \frac{\partial T_\mu z}{\partial x} \right\|^2 \geq 0, \end{aligned}$$

从而,对几乎所有的 $t \in (0, T)$,

$$\lim_{\mu \rightarrow 0} \mu T_\mu z(\cdot, t) = 0, \lim_{\mu \rightarrow 0} \frac{\partial T_\mu z(\cdot, t)}{\partial x} = 0,$$

$$\lim_{\mu \rightarrow 0} \frac{\partial^2 T_\mu z(\cdot, t)}{\partial x^2} = 0, \lim_{\mu \rightarrow 0} \frac{\partial^3 T_\mu z(\cdot, t)}{\partial x^3} = 0.$$

因此,对任意的函数 $\varphi \in H_0^2(Q_T)$,在区域 Q_T 上对方程

$$\left(\frac{\partial^4 T_\mu z}{\partial x^4} + \mu T_\mu z\right)\varphi = z\varphi$$

进行积分.根据分部积分法,有

$$\iint_{Q_T} \left(\frac{\partial^2 T_\mu z}{\partial x^2} \frac{\partial \varphi}{\partial x^2} + \mu T_\mu z \cdot \varphi\right) dxdt = \iint_{Q_T} z\varphi dxdt,$$

因此,可以得到

$$\begin{aligned} & \left(\iint_{Q_T} z\varphi dxdt\right)^2 \leq C_\varphi \iint_{Q_T} \left(\frac{\partial^2 T_\mu z}{\partial x^2}\right)^2 dxdt + \\ & C_\varphi \mu \iint_{Q_T} (T_\mu z)^2 dxdt \rightarrow 0, (\mu \rightarrow 0). \end{aligned}$$

这表明对几乎处处的 Q_T 有 $z = u_1 - u_2 = 0$.

为了证明(8)式成立.对任意的函数 $\Psi(t) \in C_0^\infty(0, T)$,在(5)式中取 $\varphi(x, t) = (T_\mu z)\Psi(t)$ 并且

使用(7)式,得到

$$\begin{aligned}
 & - \iint_{Q_T} (I - \lambda\Delta)z \frac{\partial}{\partial t} ((T_{\mu}z)\Psi(t)) dxdt + \\
 & \iint_{Q_T} v \frac{\partial^4}{\partial x^4} ((T_{\mu}z)\Psi(t)) dxdt = - \iint_{Q_T} (I - \\
 & \lambda\Delta)z \frac{\partial T_{\mu}z}{\partial t} \Psi(t) dxdt - \iint_{Q_T} (I - \lambda\Delta)z T_{\mu}z \Psi'(t) \cdot \\
 & dxdt + \iint_{Q_T} v \frac{\partial^4 T_{\mu}z}{\partial x^4} \Psi(t) dxdt = - \iint_{Q_T} (I - \\
 & \lambda\Delta) \frac{\partial T_{\mu}z}{\partial t} z \Psi(t) dxdt - \iint_{Q_T} (I - \lambda\Delta)z T_{\mu}z \Psi'(t) \cdot \\
 & dxdt + \iint_{Q_T} v(z - \mu T_{\mu}z) \Psi(t) dxdt = \iint_{Q_T} (v - \\
 & \mu T_{\mu}v) z \Psi(t) dxdt - \iint_{Q_T} (I - \lambda\Delta)z T_{\mu}z \Psi'(t) dxdt + \\
 & \iint_{Q_T} z(v - \mu T_{\mu}v) \Psi(t) dxdt = - \iint_{Q_T} (I - \\
 & \lambda\Delta)z T_{\mu}z \Psi'(t) dxdt + 2 \iint_{Q_T} z(v - \mu T_{\mu}v) \Psi(t) dxdt = \\
 & 0,
 \end{aligned}$$

即

$$\begin{aligned}
 & - \int_0^T \Psi'(t) dt \int_0^1 (I - \lambda\Delta)z T_{\mu}z dx + \int_0^T \Psi(t) dt \cdot \\
 & 2 \int_0^1 z(v - \mu T_{\mu}v) dx = - \int_0^T \Psi'(t) g_{\mu}(t) dt + \\
 & \int_0^T \Psi(t) dt 2 \int_0^1 z(v - \mu T_{\mu}v) dx = 0.
 \end{aligned}$$

因此,

$$g'_{\mu}(t) = 2 \int_0^1 (\mu T_{\mu}v - v) z dx, a. e. t \in (0, T). \tag{9}$$

由此知道 $g'_{\mu}(t) \in L^1(0, T)$ 并且 $g_{\mu}(t)$ 在 $[0, T]$ 上是一个绝对连续函数. 记

$$\Psi_{\epsilon}(t) = \int_t^{\infty} \alpha_{\epsilon}(s - \epsilon) ds,$$

其中 $\alpha_{\epsilon}(s)$ 是一维磨光核. 在(5)式中取 $\varphi(x, t) = (T_{\mu}z)\Psi_{\epsilon}(t)$, 那么

$$\begin{aligned}
 & - \iint_{Q_T} (I - \lambda\Delta)z \frac{\partial}{\partial t} ((T_{\mu}z)\Psi_{\epsilon}(t)) dxdt + \\
 & \iint_{Q_T} v \frac{\partial^4}{\partial x^4} ((T_{\mu}z)\Psi_{\epsilon}(t)) dxdt = \iint_{Q_T} (I -
 \end{aligned}$$

$$\lambda\Delta)z T_{\mu}z \alpha_{\epsilon}(t - \epsilon) dxdt + \iint_{Q_T} \Psi_{\epsilon}(t) v \frac{\partial^4 T_{\mu}z}{\partial x^4} dxdt = 0.$$

因此, 根据控制收敛定理^[4], 得到

$$\begin{aligned}
 g_{\mu}(0) &= \lim_{\epsilon \rightarrow 0} \int_0^{2\epsilon} \alpha_{\epsilon}(t - \epsilon) g_{\mu}(t) dx = \lim_{\epsilon \rightarrow 0} \iint_{Q_T} \alpha_{\epsilon}(t - \\
 & \epsilon) (I - \lambda\Delta)z T_{\mu}z dxdt = \\
 & - \lim_{\epsilon \rightarrow 0} \iint_{Q_T} \Psi_{\epsilon}(t) v \frac{\partial^4 T_{\mu}z}{\partial x^4} dxdt = 0. \tag{10}
 \end{aligned}$$

结合(6)式、(9)式和(10)式, 并注意 $\Phi(s)$ 是非减函数以及 z, v 有相同的符号, 有

$$\begin{aligned}
 0 \leq g_{\mu}(t) &= g_{\mu}(t) - g_{\mu}(0) = \int_0^t g'_{\mu}(s) ds = \\
 & \int_0^t ds \cdot 2 \int_0^1 (\mu T_{\mu}v - v) z dx = \\
 & \int_0^t ds \cdot 2 \int_0^1 (\mu T_{\mu}v) z dx - \int_0^t ds \cdot 2 \int_0^1 v z dx \leq 2 \int_0^t ds \cdot \\
 & \int_0^1 (\mu T_{\mu}v) z dx \leq \\
 & \sqrt{\mu} \iint_{Q_T} \mu (T_{\mu}v)^2 dxdt + \sqrt{\mu} \iint_{Q_T} z^2 dxdt \leq \\
 & \sqrt{\mu} \iint_{Q_T} (v^2 + z^2) dxdt \rightarrow 0, (\mu \rightarrow 0).
 \end{aligned}$$

定理 1 证明完毕.

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